

Upscaling: perforated domains

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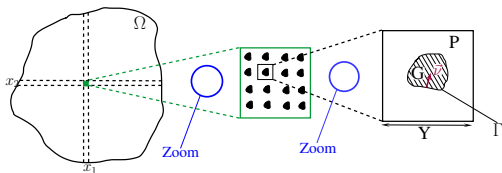
UHASSELT

KNOWLEDGE IN ACTION

The CMAT logo is set against a colorful, abstract background with a gradient from yellow to red and blue. The text "CMAT" is in white, bold, sans-serif font, centered within a dark blue, horizontally-oriented oval shape.

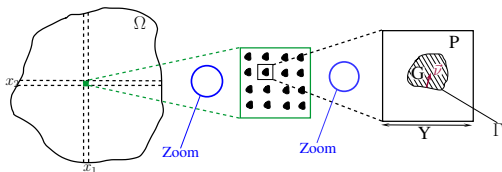
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Perforated domains



Porous medium (macro-scale, Darcy): $\Omega \subset \mathbb{R}^2$, (or $\Omega \subset \mathbb{R}^d$, $d \geq 1$)

Perforated domains



Porous medium (macro-scale, Darcy): $\Omega \subset \mathbb{R}^2$, (or $\Omega \subset \mathbb{R}^d$, $d \geq 1$)

Pore space (micro-scale, pores): $\Omega^\varepsilon \subset \Omega$ ($0 < \varepsilon \ll 1$)

With $Y = [0, 1]^2$ (**unit cell**), decompose $Y = \mathcal{P} \cup \mathcal{G} \cup \Gamma$ (*pore, grain, wall*);

For $k = (k_1, k_2)^T \in \mathbb{Z}^2$, define $\Theta_k = \{x + k/x \in \Theta\}$, and

$$\varepsilon\Theta = \{\varepsilon x/x \in \Theta\} \text{ (with } \Theta \in \{\mathcal{P}, \mathcal{G}, \Gamma\}\text{)}$$

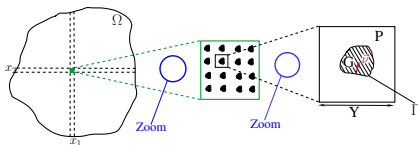
Assume $\bar{\Omega} = \cup \{\varepsilon(k + \bar{Y})/k \in \mathcal{K}_\varepsilon\}$, with \mathcal{K}_ε - set of indices;

Define

$$\Omega^\varepsilon = \cup \{\varepsilon(k + \mathcal{P})/k \in \mathcal{K}_\varepsilon\} \text{ (the total pore/void space)}$$

$$\Gamma^\varepsilon = \cup \{\varepsilon(k + \Gamma)/k \in \mathcal{K}_\varepsilon\} \text{ (the grain boundaries/pore walls)}$$

Main ideas



1. Two scales:

$$x \longrightarrow (x, y), \text{ with } y = \frac{x}{\varepsilon}$$

2. Homogenization ansatz:

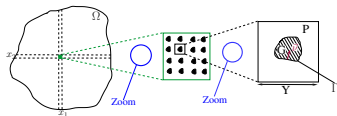
$$u^\varepsilon(x) = u_0(x, y) + \varepsilon u_1(x, y) + \varepsilon^2 u_2(x, y) + \dots,$$

with u_k being Y -periodic (w.r.t. y) ($k = 0, 1, \dots$).

3. Derivatives: Consider $g^\varepsilon(x) := g(x, \frac{x}{\varepsilon}) = g(x, y)$; then

$$\frac{\partial g}{\partial x_i}(x, y) \quad \text{becomes} \quad \frac{dg^\varepsilon}{dx_i}(x) = \frac{\partial g}{\partial x_i}(x, y) + \frac{\partial y_i}{\partial x_i} \frac{\partial g}{\partial y_i}(x, y) = \frac{\partial g}{\partial x_i}(x, y) + \frac{1}{\varepsilon} \frac{\partial g}{\partial y_i}(x, y)$$

Periodicity and micro-scale model

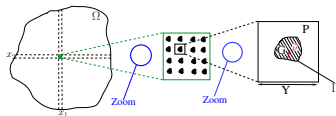


Let $A^\varepsilon : \mathbb{R}^2 \rightarrow \mathbb{R}$ - rapidly oscillating scalar function (can be extended to matrices) - defined by $A^\varepsilon(x) = A(x, \frac{x}{\varepsilon})$ for all $x = (x_1, x_2)^T \in \mathbb{R}^2$, where $A : \mathbb{R}^4 \rightarrow \mathbb{R}$ satisfies for all $x = (x_1, x_2)^T$ and $y = (y_1, y_2)^T$

Bounded : $0 < m \leq A(x_1, x_2, y_1, y_2) \leq M < \infty$,

Periodic : $A(x_1, x_2, y_1, y_2) = A(x_1, x_2, y_1 + 1, y_2)$
 $= A(x_1, x_2, y_1, y_2 + 1) = A(x_1, x_2, y_1 + 1, y_2 + 1) = \dots$

Periodicity and micro-scale model



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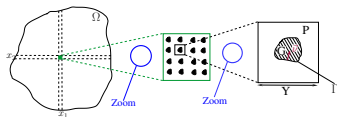
$$\begin{aligned} \text{Periodic : } \quad A(x_1, x_2, y_1, y_2) &= A(x_1, x_2, y_1 + 1, y_2) \\ &= A(x_1, x_2, y_1, y_2 + 1) = A(x_1, x_2, y_1 + 1, y_2 + 1) = \dots \end{aligned}$$

Consider the diffusion problem in the perforated domain (porous medium)

$$(P^\varepsilon) \quad \begin{cases} -\nabla \cdot (A^\varepsilon(x) \nabla u^\varepsilon) &= f, & \text{for all } x \in \Omega^\varepsilon \text{ (pore space),} \\ -\nu \cdot (A^\varepsilon(x) \nabla u^\varepsilon) &= 0, & \text{at } \Gamma^\varepsilon \text{ (pore walls),} \\ u^\varepsilon &= 0, & \text{at } \partial\Omega \text{ (outer boundary).} \end{cases}$$

Note: The equation is defined only in the pore space Ω^ε , not in the grain space $\Omega \setminus \bar{\Omega}^\varepsilon$!

Asymptotic expansions

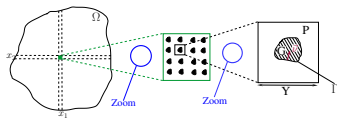


As before, take $y = \frac{x}{\varepsilon}$ and use

$u^\varepsilon(x) = u_0(x, y) + \varepsilon u_1(x, y) + \varepsilon^2 u_2(x, y) + \dots$, with u_k being Y -periodic;

$A^\varepsilon(x) = A(x, y)$ and $\nabla \rightarrow \nabla_x + \frac{1}{\varepsilon} \nabla_y$

Asymptotic expansions



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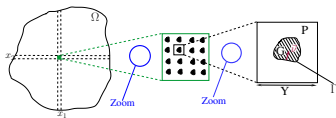
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$A^\varepsilon(x) = A(x, y)$ and $\nabla \rightarrow \nabla_x + \frac{1}{\varepsilon} \nabla_y$

Equation (in Ω^ε)

$$\begin{aligned} & -\frac{1}{\varepsilon^2} \nabla_y \cdot (A(x, y) \nabla_y u_0(x, y)) \\ & -\frac{1}{\varepsilon} \{ \nabla_x \cdot [A(x, y) \nabla_y u_0(x, y)] + \nabla_y \cdot [A(x, y) (\nabla_x u_0(x, y) + \nabla_y u_1(x, y))] \} \\ & - \{ \nabla_x \cdot [A(x, y) (\nabla_x u_0(x, y) + \nabla_y u_1(x, y))] + \nabla_y \cdot [A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] \} \\ & = f + O(\varepsilon), \quad \text{in } \Omega^\varepsilon. \end{aligned}$$

Asymptotic expansions



As before, take $y = \frac{x}{\varepsilon}$ and use

$$u^\varepsilon(x) = u_0(x, y) + \varepsilon u_1(x, y) + \varepsilon^2 u_2(x, y) + \dots, \text{ with } u_k \text{ being } Y\text{-periodic};$$

$$A^\varepsilon(x) = A(x, y) \quad \text{and} \quad \nabla \longrightarrow \nabla_x + \frac{1}{\varepsilon} \nabla_y$$

Equation (in Ω^ε)

$$\begin{aligned} & -\frac{1}{\varepsilon^2} \nabla_y \cdot (A(x, y) \nabla_y u_0(x, y)) \\ & -\frac{1}{\varepsilon} \{ \nabla_x \cdot [A(x, y) \nabla_y u_0(x, y)] + \nabla_y \cdot [A(x, y) (\nabla_x u_0(x, y) + \nabla_y u_1(x, y))] \} \\ & - \{ \nabla_x \cdot [A(x, y) (\nabla_x u_0(x, y) + \nabla_y u_1(x, y))] + \nabla_y \cdot [A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] \} \\ & = f + O(\varepsilon), \quad \text{in } \Omega^\varepsilon. \end{aligned}$$

Boundary condition (at Γ^ε)

$$\begin{aligned} 0 = & -\frac{1}{\varepsilon} \{ \nu \cdot (A(x, y) \nabla_y u_0(x, y)) \} - \{ \nu \cdot (A(x, y) (\nabla_x u_0(x, y) + \nabla_y u_1(x, y))) \} \\ & - \varepsilon \{ \nu \cdot (A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))) \} + O(\varepsilon^2), \quad \text{on } \Gamma^\varepsilon. \end{aligned}$$

Lowest order terms

$$\begin{aligned} & -\frac{1}{\varepsilon^2} \nabla_y \cdot (A(x, y) \nabla_y u_0(x, y)) \\ & -\frac{1}{\varepsilon} \{ \nabla_x \cdot [A(x, y) \nabla_y u_0(x, y)] + \nabla_y \cdot [A(x, y) (\nabla_x u_0(x, y) + \nabla_y u_1(x, y))] \} \\ & - \{ \nabla_x \cdot [A(x, y) (\nabla_x u_0(x, y) + \nabla_y u_1(x, y))] + \nabla_y \cdot [A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] \} \\ = & f + O(\varepsilon), \quad \text{in } \Omega^\varepsilon. \\ 0 = & -\frac{1}{\varepsilon} \{ \nu \cdot (A(x, y) \nabla_y u_0(x, y)) \} - \{ \nu \cdot (A(x, y) (\nabla_x u_0(x, y) + \nabla_y u_1(x, y))) \} \\ & - \varepsilon \{ \nu \cdot (A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))) \} + O(\varepsilon^2), \quad \text{on } \Gamma^\varepsilon. \end{aligned}$$

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$$\text{(Problem } P^{-2}) \quad \begin{cases} -\nabla_y \cdot (A(x, y) \nabla_y u_0(x, y)) = 0, & \text{for all } y \in \mathcal{P}, \\ -\nu \cdot (A(x, y) \nabla_y u_0(x, y)) = 0, & \text{for all } y \in \Gamma, \\ u_0(x, y) \text{ is } Y\text{-periodic.} \end{cases}$$

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Note: The only possible solution is y -independent, $u_0(x, y) = u_0(x)$, so $\nabla_y u_0(x, y) = 0$ for all $y \in \mathcal{P}$. Again, x plays the role of a parameter.

Q: $u_0(x) = ?$

Correctors

$$\begin{aligned} & -\frac{1}{\varepsilon^2} \nabla_y \cdot (A(x, y) \nabla_y u_0(x)) \\ & -\frac{1}{\varepsilon} \{ \nabla_x \cdot [A(x, y) \nabla_y u_0(x)] + \nabla_y \cdot [A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))] \} \\ & - \{ \nabla_x \cdot [A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))] + \nabla_y \cdot [A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] \} \\ = & f + O(\varepsilon), \quad \text{in } \Omega^\varepsilon. \\ 0 = & -\frac{1}{\varepsilon} \{ \nu \cdot (A(x, y) \nabla_y u_0(x)) \} - \{ \nu \cdot (A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))) \} \\ & - \varepsilon \{ \nu \cdot (A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))) \} + O(\varepsilon^2), \quad \text{on } \Gamma^\varepsilon. \end{aligned}$$

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$$(Problem P^{-1}) \quad \begin{cases} -\nabla_y \cdot [A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))] = 0, & \text{for all } y \in \mathcal{P}, \\ -\nu \cdot (A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))) = 0, & \text{for all } y \in \Gamma, \\ u_1(x, y) \text{ is } Y\text{-periodic.} \end{cases}$$

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Since $\nabla_x u_0(x) = \sum_{j=1}^2 e_j \partial_{x_j} u_0(x)$, replace $\nabla_x u_0(x)$ by e_j and consider again the *cell problems*

$$(P_j^{-1})(x) \quad \begin{cases} -\nabla_y \cdot (A(x, y) \nabla_y w^j(x, y)) = \nabla_y \cdot (A(x, y) e_j), & \text{for all } y \in \mathcal{P}, \\ -\nu \cdot (A(x, y) \nabla_y w^j(x, y)) = \nu \cdot (A(x, y) e_j), & \text{for all } y \in \Gamma, \\ w^j(x, y) \text{ is } Y\text{-periodic,} & \int_{\mathcal{P}} w^j(x, y) dy = 0. \end{cases}$$

Correctors

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As before, this gives

$$u_1(x, y) = \tilde{u}_1(x) + \sum_{j=1}^2 w^j(x, y) \partial_{x_j} u_0(x),$$

Note: Now w^j depends explicitly on x !

Perforated domains: 0-order terms

$$\begin{aligned} & -\frac{1}{\varepsilon^2} \nabla_y \cdot (A(x, y) \nabla_y u_0(x)) \\ & -\frac{1}{\varepsilon} \{ \nabla_x \cdot [A(x, y) \nabla_y u_0(x)] + \nabla_y \cdot [A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))] \} \\ & - \{ \nabla_x \cdot [A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))] + \nabla_y \cdot [A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] \} \\ = & \mathbf{f} + O(\varepsilon), \quad \text{in } \Omega^\varepsilon. \\ 0 = & -\frac{1}{\varepsilon} \{ \nu \cdot (A(x, y) \nabla_y u_0(x)) \} - \{ \nu \cdot (A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))) \} \\ & - \varepsilon \{ \nu \cdot (A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))) \} + O(\varepsilon^2), \quad \text{on } \Gamma^\varepsilon. \end{aligned}$$

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 0 = & -\frac{1}{\varepsilon} \{ \nu \cdot (A(x, y) \nabla_y u_0(x)) \} - \{ \nu \cdot (A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))) \} \\
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 \end{aligned}$$

$$(P^0)(x) \left\{ \begin{array}{ll}
 -\nabla_x \cdot [A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))] & \\
 -\nabla_y \cdot [A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] & = f, \text{ for all } y \in \mathcal{P}, \\
 -\nu \cdot (A(x, y) (\nabla_y u_2(x, y) + \nabla_x u_1(x, y))) & = 0 \text{ for all } y \in \Gamma, \\
 u_2(x, y) & \text{is } Y\text{-periodic.}
 \end{array} \right.$$

Perforated domains: 0-order terms

$$\begin{aligned} & -\frac{1}{\varepsilon^2} \nabla_y \cdot (A(x, y) \nabla_y u_0(x)) \\ & -\frac{1}{\varepsilon} \{ \nabla_x \cdot [A(x, y) \nabla_y u_0(x)] + \nabla_y \cdot [A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))] \} \\ & - \{ \nabla_x \cdot [A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))] + \nabla_y \cdot [A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] \} \\ = & \mathbf{f} + O(\varepsilon), \quad \text{in } \Omega^\varepsilon. \end{aligned}$$

$$\begin{aligned} 0 = & -\frac{1}{\varepsilon} \{ \nu \cdot (A(x, y) \nabla_y u_0(x)) \} - \{ \nu \cdot (A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))) \} \\ & - \varepsilon \{ \nu \cdot (A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))) \} + O(\varepsilon^2), \quad \text{on } \Gamma^\varepsilon. \end{aligned}$$

$$(P^0)(x) \left\{ \begin{array}{l} -\nabla_x \cdot [A(x, y) (\nabla_x u_0(x) + \nabla_y u_1(x, y))] \\ -\nabla_y \cdot [A(x, y) (\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] \\ -\nu \cdot (A(x, y) (\nabla_y u_2(x, y) + \nabla_x u_1(x, y))) \\ u_2(x, y) \end{array} \right. \begin{array}{l} = f, \text{ for all } y \in \mathcal{P}, \\ = 0 \text{ for all } y \in \Gamma, \\ \text{is } Y\text{-periodic.} \end{array}$$

Note: Again, we have three unknown functions u_0 , u_1 and u_2 . We know how the first two are related, but nothing about u_2 .

Perforated domains: 0-order terms

$$(P^0)(x) \begin{cases} -\nabla_x \cdot [A(x, y)(\nabla_x u_0(x) + \nabla_y u_1(x, y))] \\ -\nabla_y \cdot [A(x, y)(\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] \\ -\nu \cdot (A(x, y)(\nabla_y u_2(x, y) + \nabla_x u_1(x, y))) \\ u_2(x, y) \end{cases} \begin{array}{l} = f, \text{ for all } y \in \mathcal{P}, \\ = 0 \text{ for all } y \in \Gamma, \\ \text{is } Y\text{-periodic.} \end{array}$$

Idea: Eliminate u_2 by integration!

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Idea: Eliminate u_2 by integration!

For all $x \in \Omega$, with $|\mathcal{P}| = \int_{\mathcal{P}} dy$ (the volume of the void space in a reference cell, the *porosity*) one has

$$\begin{aligned} & -\nabla_x \cdot \left[\int_{\mathcal{P}} A(x, y)(\nabla_x u_0(x) + \nabla_y u_1(x, y)) dy \right] \\ & - \int_{\mathcal{P}} \nabla_y \cdot [A(x, y)(\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] dy = |\mathcal{P}| f(x). \end{aligned}$$

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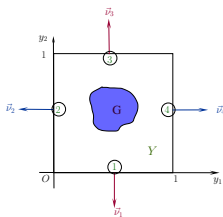
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Note:

$$\begin{aligned} & \int_{\mathcal{P}} \nabla_y \cdot [A(x, y)(\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] dy \\ & = \int_{\partial \mathcal{P}} \nu \cdot [A(x, y)(\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] d\sigma_y = 0. \end{aligned}$$



Perforated domains: upscaled/homogenized equation

$$(P^0)(x) \left\{ \begin{array}{ll} -\nabla_x \cdot [A(x, y)(\nabla_x u_0(x) + \nabla_y u_1(x, y))] & \\ -\nabla_y \cdot [A(x, y)(\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] & = f, \text{ for all } y \in \mathcal{P}, \\ -\nu \cdot (A(x, y)(\nabla_y u_2(x, y) + \nabla_x u_1(x, y))) & = 0 \text{ for all } y \in \Gamma, \\ u_2(x, y) & \text{is } Y\text{-periodic.} \end{array} \right.$$

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With $u_1(x, y) = \tilde{u}_1(x) + \sum_{j=1}^2 w^j(x, y) \partial_{x_j} u_0(x)$, this gives

$$-\nabla_x \cdot \left[\int_{\mathcal{P}} A(x, y) \left(\nabla_x u_0(x) + \sum_{j=1}^d \partial_{x_j} u_0(x) \nabla_y w^j(x, y) \right) dy \right] = |\mathcal{P}| f(x).$$

Perforated domains: upscaled/homogenized equation

$$(P^0)(x) \begin{cases} -\nabla_x \cdot [A(x, y)(\nabla_x u_0(x) + \nabla_y u_1(x, y))] \\ \quad -\nabla_y \cdot [A(x, y)(\nabla_x u_1(x, y) + \nabla_y u_2(x, y))] & = f, \text{ for all } y \in \mathcal{P}, \\ \quad -\nu \cdot (A(x, y)(\nabla_y u_2(x, y) + \nabla_x u_1(x, y))) & = 0 \text{ for all } y \in \Gamma, \\ u_2(x, y) & \text{is } Y\text{-periodic.} \end{cases}$$

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Upscaled model (after dividing by $|\mathcal{P}|$, and with $U = u_0(x)$):

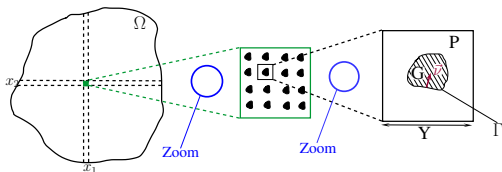
$$(P) \quad \begin{cases} -\nabla \cdot (A^*(x) \nabla U) & = f, \quad \text{for all } x \in \Omega, \\ U & = 0, \quad \text{on } \partial\Omega, \end{cases}$$

with $A^* : \Omega \rightarrow \mathbb{R}^{d \times d}$ having the elements $a_{ij}^* : \Omega \rightarrow \mathbb{R}$, $(i, j \in \{1, 2\})$

$$a_{ij}^* : \Omega \rightarrow \mathbb{R}, \quad (i, j \in \{1, 2\}) a_{ij}^*(x) = \frac{1}{|\mathcal{P}|} \int_{\mathcal{P}} (A(x, y)(e_j + \nabla_y w^j(x, y))) \cdot e_i dy.$$

Note: A^* is symmetric and positive definite

The derivation of the Darcy law



Porous medium (macro-scale, Darcy): $\Omega \subset \mathbb{R}^2$, (or $\Omega \subset \mathbb{R}^d$, $d \geq 1$)

Pore space (micro-scale, pores): $\Omega^\varepsilon \subset \Omega$ ($0 < \varepsilon \ll 1$)

With $Y = [0, 1]^2$ (unit cell), decompose $Y = \mathcal{P} \cup \mathcal{G} \cup \Gamma$ (pore, grain, wall);

For $k = (k_1, k_2)^T \in \mathbb{Z}^2$, define $\Theta_k = \{x + k/x \in \Theta\}$, and

$$\varepsilon\Theta = \{\varepsilon x/x \in \Theta\} \text{ (with } \Theta \in \{\mathcal{P}, \mathcal{G}, \Gamma\}\text{)}$$

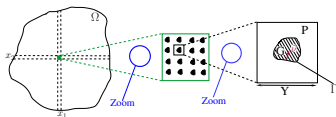
Assume $\bar{\Omega} = \cup \{\varepsilon(k + \bar{Y})/k \in \mathcal{K}_\varepsilon\}$, with \mathcal{K}_ε - set of indices;

Define

$$\Omega^\varepsilon = \cup \{\varepsilon(k + \mathcal{P})/k \in \mathcal{K}_\varepsilon\} \text{ (the total pore/void space)}$$

$$\Gamma^\varepsilon = \cup \{\varepsilon(k + \Gamma)/k \in \mathcal{K}_\varepsilon\} \text{ (the grain boundaries/pore walls)}$$

The micro-scale flow model

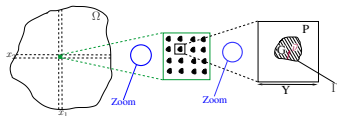


Stokes flow at the scale of pores:

$$\begin{cases} \varepsilon^2 \mu \Delta \mathbf{q}^\varepsilon = \nabla p^\varepsilon(x), & \text{in } \Omega^\varepsilon \\ \nabla \cdot \mathbf{q}^\varepsilon(x) = 0, & \text{in } \Omega^\varepsilon \\ \mathbf{q}^\varepsilon(x) = 0, & \text{on } \Gamma^\varepsilon, \end{cases}$$

Note: Small viscosity, $O(\varepsilon^2)$, or large pressure gradients

Asymptotic expansion



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Note: Small viscosity, $O(\varepsilon^2)$, or large pressure gradients

Homogenization ansatz: with x replaced by $(x, y = x/\varepsilon)$

$$\mathbf{q}^\varepsilon = \mathbf{q}_0(x, y) + \varepsilon \mathbf{q}_1(x, y) + \varepsilon^2 \mathbf{q}_2(x, y) + \dots$$

$$p^\varepsilon = p_0(x, y) + \varepsilon p_1(x, y) + \varepsilon^2 p_2(x, y) + \dots$$

$$\nabla \rightarrow \nabla_x + \frac{1}{\varepsilon} \nabla_y, \quad \Delta \rightarrow \Delta_x + \frac{1}{\varepsilon} \nabla_x \cdot \nabla_y + \frac{1}{\varepsilon} \nabla_y \cdot \nabla_x + \frac{1}{\varepsilon^2} \Delta_y$$

with \mathbf{q}_k, p_k - Y -periodic.

Asymptotic expansion

Homogenization ansatz: with x replaced by $(x, y = x/\varepsilon)$

$$\mathbf{q}^\varepsilon = \mathbf{q}_0(x, y) + \varepsilon \mathbf{q}_1(x, y) + \varepsilon^2 \mathbf{q}_2(x, y) + \dots$$

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with $\mathbf{q}_k, \mathbf{p}_k$ - Y -periodic.

$$\left\{ \begin{array}{ll} \mu \Delta_y \mathbf{q}_0(x, y) + \varepsilon \mu [(\nabla_x \cdot \nabla_y + \nabla_y \cdot \nabla_x) \mathbf{q}_0 + \Delta_y \mathbf{q}_1] + O(\varepsilon^2) \\ = \frac{1}{\varepsilon} \nabla_y \mathbf{p}_0 + (\nabla_x \mathbf{p}_0 + \nabla_y \mathbf{p}_1) + \varepsilon (\nabla_x \mathbf{p}_1 + \nabla_y \mathbf{p}_2) + O(\varepsilon^2), & \text{in } \Omega^\varepsilon \\ \frac{1}{\varepsilon} \nabla_y \cdot \mathbf{q}_0 + (\nabla_x \cdot \mathbf{q}_0 + \nabla_y \cdot \mathbf{q}_1) + \varepsilon (\nabla_x \cdot \mathbf{q}_1 + \nabla_y \cdot \mathbf{q}_2) + O(\varepsilon^2) = 0, & \text{in } \Omega^\varepsilon \\ \mathbf{q}_0 + \varepsilon \mathbf{q}_1 + \varepsilon^2 \mathbf{q}_2 + O(\varepsilon^2) = 0, & \text{on } \Gamma^\varepsilon. \end{array} \right.$$

The derivation of the Darcy law: y -independent pressure

$$\left\{ \begin{array}{ll} \mu \Delta_y q_0(x, y) + \varepsilon \mu [(\nabla_x \cdot \nabla_y + \nabla_y \cdot \nabla_x) q_0 + \Delta_y q_1] + O(\varepsilon^2) \\ \quad = \frac{1}{\varepsilon} \nabla_y p_0 + (\nabla_x p_0 + \nabla_y p_1) + \varepsilon (\nabla_x p_1 + \nabla_y p_2) + O(\varepsilon^2), & \text{in } \Omega^\varepsilon \\ \frac{1}{\varepsilon} \nabla_y \cdot q_0 + (\nabla_x \cdot q_0 + \nabla_y \cdot q_1) + \varepsilon (\nabla_x \cdot q_1 + \nabla_y \cdot q_2) + O(\varepsilon^2) = 0, & \text{in } \Omega^\varepsilon \\ q_0 + \varepsilon q_1 + \varepsilon^2 q_2 + O(\varepsilon^2) = 0, & \text{on } \Gamma^\varepsilon. \end{array} \right.$$

Lowest order term:

Problem P^{-1} $\nabla_y p_0 = 0$, for all $y \in \mathcal{P}$

gives $p_0 = p_0(x)$.

The derivation of the Darcy law: 0-order velocity

$$\left\{ \begin{array}{ll} \mu \Delta_y \mathbf{q}_0(x, y) + \varepsilon \mu [(\nabla_x \cdot \nabla_y + \nabla_y \cdot \nabla_x) \mathbf{q}_0 + \Delta_y \mathbf{q}_1] + O(\varepsilon^2) \\ \quad = \frac{1}{\varepsilon} \nabla_y p_0 + (\nabla_x p_0 + \nabla_y p_1) + \varepsilon (\nabla_x p_1 + \nabla_y p_2) + O(\varepsilon^2), & \text{in } \Omega^\varepsilon \\ \frac{1}{\varepsilon} \nabla_y \cdot \mathbf{q}_0 + (\nabla_x \cdot \mathbf{q}_0 + \nabla_y \cdot \mathbf{q}_1) + \varepsilon (\nabla_x \cdot \mathbf{q}_1 + \nabla_y \cdot \mathbf{q}_2) + O(\varepsilon^2) = 0, & \text{in } \Omega^\varepsilon \\ \mathbf{q}_0 + \varepsilon \mathbf{q}_1 + \varepsilon^2 \mathbf{q}_2 + O(\varepsilon^2) = 0, & \text{on } \Gamma^\varepsilon. \end{array} \right.$$

Next terms (w.r.t. y):

$$\text{Problem } P^0 \left\{ \begin{array}{ll} \mu \Delta_y \mathbf{q}_0(x, y) = \nabla_x p_0(x) + \nabla_y p_1(x, y), & \text{in } \mathcal{P} \\ \nabla_y \cdot \mathbf{q}_0(x, y) = 0, & \text{in } \mathcal{P} \\ \mathbf{q}_0(x, y) = 0, & \text{on } \Gamma \\ \mathbf{q}_0, p_1 - Y\text{-periodic.} \end{array} \right.$$

The derivation of the Darcy law: cell problems

$$\text{Problem } P^0 \quad \left\{ \begin{array}{ll} \mu \Delta_y q_0(x, y) = \nabla_x p_0(x) + \nabla_y p_1(x, y), & \text{in } \mathcal{P} \\ \nabla_y \cdot q_0(x, y) = 0, & \text{in } \mathcal{P} \\ q_0(x, y) = 0, & \text{on } \Gamma \\ q_0, p_1 - Y\text{-periodic.} & \end{array} \right.$$

Idea: Express the p_0 -dependence of q_0 through “cell problems”.

Since $\nabla_x p_0(x) = \sum_{j=1}^d e_j \partial_{x_j} p_0(x)$, solve

$$(P_j) \quad \left\{ \begin{array}{ll} -\Delta_y \chi^j(y) = \nabla_y \Pi^j(y) + e_j, & \text{in } \mathcal{P} \\ \nabla_y \cdot \chi^j(y) = 0, & \text{in } \mathcal{P} \\ \chi^j(y) = 0, & \text{on } \Gamma \\ \chi^j, \Pi^j - Y\text{-periodic,} & \end{array} \right.$$

and identify

$$q_0(x, y) = -\frac{1}{\mu} \sum_{j=1}^d \partial_{x_j} p_0(x) \cdot \chi^j(y), \quad p_1(x, y) = \sum_{j=1}^d \partial_{x_j} p_0(x) \cdot \Pi^j(y).$$

The Darcy law (upscaled model)

The *averaged velocity*

$$\mathbf{q}(x) = \frac{1}{|\mathcal{P}|} \int_{\mathcal{P}} \mathbf{q}_0(x, y) dy,$$

satisfies the Darcy law

$$\mathbf{q}(x) = -\frac{1}{\mu} \underline{\mathbf{K}} \nabla p_0(x) \quad (\text{for all } x \text{ in } \Omega),$$

where the *permeability tensor* $\underline{\mathbf{K}}$ is defined by $(\chi^j = (\chi_1^j, \dots, \chi_d^j)^T)$

$$k_{ij} = \frac{1}{|\mathcal{P}|} \int_{\mathcal{P}} \chi_i^j(y) dy, \quad \text{for all } i, j = 1, \dots, d.$$

The Darcy law: properties

The *averaged velocity*

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satisfies the Darcy law

$$\mathbf{q}(x) = -\frac{1}{\mu} \underline{\mathbf{K}} \nabla p_0(x) \quad (\text{for all } x \text{ in } \Omega),$$

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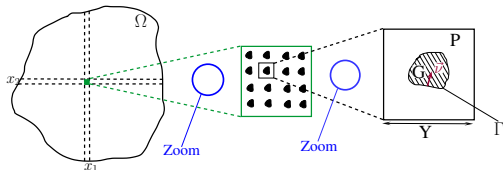
$$k_{ij} = \frac{1}{|\mathcal{P}|} \int_{\mathcal{P}} \chi_i^j(y) dy, \quad \text{for all } i, j = 1, \dots, d.$$

Note: Since $\nabla_x \cdot \mathbf{q}_0 + \nabla_y \cdot \mathbf{q}_1 = 0$, \mathbf{q} is divergence-free!

$$\begin{aligned} \nabla_x \cdot \mathbf{q}(x) &= \frac{1}{|\mathcal{P}|} \int_{\mathcal{P}} \nabla_x \cdot \mathbf{q}_0(x, y) dy \\ &= -\frac{1}{|\mathcal{P}|} \int_{\mathcal{P}} \nabla_y \cdot \mathbf{q}_1(x, y) dy \\ &= -\frac{1}{|\mathcal{P}|} \int_{\partial \mathcal{P}} \nu \cdot \mathbf{q}_1(x, y) d\sigma_y \\ &= -\frac{1}{|\mathcal{P}|} \int_{\partial \mathcal{Y}} \nu \cdot \mathbf{q}_1(x, y) d\sigma_y - \frac{1}{|\mathcal{P}|} \int_{\Gamma} \nu \cdot \mathbf{q}_1(x, y) d\sigma_y = 0. \end{aligned}$$

Lemma. The tensor $\underline{\mathbf{K}}$ is symmetric and positive definite.

The pore-scale reactive flow model



In Ω_P^ε :

$$\begin{cases} \varepsilon^2 \mu \Delta q = \nabla p, \\ \nabla \cdot q = 0, \\ \partial_t u = D \Delta u - \nabla \cdot (qu). \end{cases}$$

At Γ_G^ε :

$$\begin{cases} q = \bar{0}, \\ -D\nu \cdot \nabla u = \varepsilon \partial_t v, \\ \partial_t v = D_a (r_S - r_I). \end{cases}$$

Rem: Non-equilibrium kinetics, $Da = O(1)$.

Pb: Find an upscaled model!

Asymptotic expansion

In Ω_p^ε :

$$\begin{cases} \varepsilon^2 \mu \Delta q &= \nabla p, \\ \nabla \cdot q &= 0, \\ \partial_t u &= D \Delta u - \nabla \cdot (qu). \end{cases}$$

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Asymptotic expansion

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$$\begin{cases} q &= \bar{0}, \\ -D\nu \cdot \nabla u &= \varepsilon \partial_t v, \\ \partial_t v &= D_a(r_S - r_I). \end{cases}$$

Fast/slow variable:

$$x \longrightarrow (x, y), \text{ with } \boxed{y = \frac{x}{\varepsilon}}, \text{ giving } \frac{\partial u}{\partial x_i}(x) \longrightarrow \frac{\partial u}{\partial x_i}(x, y) + \frac{1}{\varepsilon} \frac{\partial u}{\partial y_i}(x, y)$$

Homogenization ansatz:

$$u^\varepsilon(t, x) = u_0(t, x, y) + \varepsilon u_1(t, x, y) + \varepsilon^2 u_2(t, x, y) + \dots,$$

$$v^\varepsilon(t, x) = v_0(t, x, y) + \varepsilon v_1(t, x, y) + \varepsilon^2 v_2(t, x, y) + \dots,$$

$$r_\alpha(u^\varepsilon, v^\varepsilon) = r_\alpha(u_0, v_0) + \varepsilon u_1 \partial_u r_\alpha(u_0, v_0) + \varepsilon u_1 \partial_u r_\alpha(u_0, v_0) + \dots, (\alpha = S, I)$$

$$q^\varepsilon(x) = q_0(x) + \varepsilon q_1(x, y) + \varepsilon^2 q_2(x, y) + \dots$$

with u_k and q_k being Y -periodic.

The reactive transport component of the model

$$\begin{aligned} \text{In } \Omega_P^\varepsilon : \quad & \partial_t u_0 + (\nabla_x + \frac{1}{\varepsilon} \nabla_y) \cdot ((q_0 + \varepsilon q_1)(u_0 + \varepsilon u_1)) \\ & = (\nabla_x + \frac{1}{\varepsilon} \nabla_y) \cdot [D (\nabla_x + \frac{1}{\varepsilon} \nabla_y) (u_0 + \varepsilon u_1 + \varepsilon^2 u_2)] + O(\varepsilon) \end{aligned}$$

$$\begin{aligned} \text{At } \Gamma_G^\varepsilon : \quad & -\nu \cdot D (\frac{1}{\varepsilon} \nabla_y u_0 + (\nabla_x u_0 + \nabla_y u_1) + \varepsilon (\nabla_x u_1 + \nabla_y u_2)) = \varepsilon \partial_t v_0 + O(\varepsilon^2) \\ & \partial_t v_0 = Da(r_S(u_0, v_0) - r_I(u_0, v_0)) + O(\varepsilon) \end{aligned}$$

The reactive transport component of the model

$$\begin{aligned} \text{In } \Omega_P^\varepsilon : \quad & \partial_t u_0 + (\nabla_x + \frac{1}{\varepsilon} \nabla_y) \cdot ((q_0 + \varepsilon q_1)(u_0 + \varepsilon u_1)) \\ & = (\nabla_x + \frac{1}{\varepsilon} \nabla_y) \cdot [D (\nabla_x + \frac{1}{\varepsilon} \nabla_y) (u_0 + \varepsilon u_1 + \varepsilon^2 u_2)] + O(\varepsilon) \end{aligned}$$

$$\begin{aligned} \text{At } \Gamma_G^\varepsilon : \quad & -\nu \cdot D \left(\frac{1}{\varepsilon} \nabla_y u_0 + (\nabla_x u_0 + \nabla_y u_1) + \varepsilon (\nabla_x u_1 + \nabla_y u_2) \right) = \varepsilon \partial_t v_0 + O(\varepsilon^2) \\ & \partial_t v_0 = Da(r_S(u_0, v_0) - r_I(u_0, v_0)) + O(\varepsilon) \end{aligned}$$

As before

$$q(x) = \frac{1}{|\mathcal{P}|} \int_{\mathcal{P}} q_0(x, y) dy = -\frac{1}{\mu} K \nabla p_0(x) \quad (\text{Darcy law})$$

$$\nabla_x \cdot q(x) = \nabla_x \cdot q_0(x, y) + \nabla_y \cdot q_1(x, y) = 0 \quad (\text{divergence free})$$

$$u_0(x, y) = u_0(x) \quad (\text{no } y\text{-dependence}), \text{ and}$$

$$u_1(x, y) = \tilde{u}_1(x) + \sum_{j=1}^d w^j(x, y) \partial_{x_j} u_0(x),$$

with w^j solving the cell problems (use $\nabla_y \cdot q_0 = 0$)

$$(P_j^{-1})(x) \begin{cases} -\nabla_y \cdot (D \nabla_y w^j(y)) & = 0, & \text{in } \mathcal{P}, \\ -\nu \cdot (D \nabla_y w^j(y)) & = \nu \cdot (D e_j), & \text{on } \Gamma, \\ w^j(y) & \text{is } Y\text{-periodic, } \int_{\mathcal{P}} w^j(y) dy = 0. \end{cases}$$

The reactive transport component: 0-order terms

The problems

$$\left\{ \begin{array}{ll} \partial_t u_0 + \nabla_x \cdot (q_0 u_0) + \nabla_y \cdot (q_0 u_1 + q_1 u_0) \\ \quad = \{ \nabla_x \cdot (D(\nabla_x u_0 + \nabla_y u_1)) + \nabla_y \cdot (D(\nabla_x u_1 + \nabla_y u_2)) \}, & \text{in } \mathcal{P} \\ \partial_t v_0 = -\nu \cdot \{ D(\nabla_x u_1 + \nabla_y u_2) \}, & \text{at } \Gamma \\ \partial_t v_0 = Da(r_S(u_0, v_0) - r_I(u_0, v_0)), & \text{at } \Gamma \end{array} \right.$$

The reactive transport component: 0-order terms

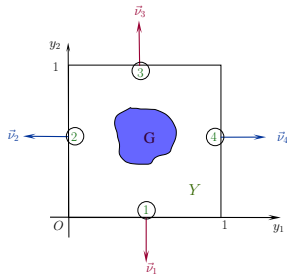
The problems

$$\left\{ \begin{array}{ll} \partial_t u_0 + \nabla_x \cdot (q_0 u_0) + \nabla_y \cdot (q_0 u_1 + q_1 u_0) \\ \quad = \{ \nabla_x \cdot (D (\nabla_x u_0 + \nabla_y u_1)) + \nabla_y \cdot (D (\nabla_x u_1 + \nabla_y u_2)) \}, & \text{in } \mathcal{P} \\ \partial_t v_0 = -\nu \cdot \{ D (\nabla_x u_1 + \nabla_y u_2) \}, & \text{at } \Gamma \\ \partial_t v_0 = Da(r_S(u_0, v_0) - r_I(u_0, v_0)), & \text{at } \Gamma \end{array} \right.$$

Integrate over \mathcal{P} (and divide by $|\mathcal{P}|$)

$$\begin{aligned} \partial_t u_0 + \nabla_x \cdot (q_0 u_0) + \frac{1}{|\mathcal{P}|} \int_{\mathcal{P}} \nabla_y \cdot (q_0 u_1 + q_1 u_0) dy \\ = \nabla_x \cdot (D \cdot \nabla_x u_0) + \frac{1}{|\mathcal{P}|} \nabla_x \cdot (\int_{\mathcal{P}} D \nabla_y u_1 dy) + \frac{1}{|\mathcal{P}|} \int_{\mathcal{P}} \nabla_y \cdot \{ D (\nabla_x u_1 + \nabla_y u_2) \} dy \end{aligned}$$

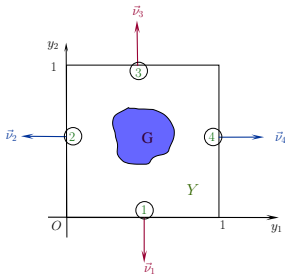
Incorporating adsorption/desorption effects



Note:

$$\int_{\mathcal{P}} \nabla_y \cdot (q_0 u_1 + q_1 u_0) dy = \int_{\partial P} \nu \cdot (q_0 u_1 + q_1 u_0) d\sigma_y$$
$$= \int_{\partial Y} \nu \cdot (q_0 u_1 + q_1 u_0) d\sigma_y + \int_{\Gamma} \nu \cdot (q_0 u_1 + q_1 u_0) d\sigma_y = 0,$$

Incorporating adsorption/desorption effects



Note:
$$\int_{\mathcal{P}} \nabla_y \cdot (q_0 u_1 + q_1 u_0) dy = \int_{\partial \mathcal{P}} \nu \cdot (q_0 u_1 + q_1 u_0) d\sigma_y$$

$$= \int_{\partial Y} \nu \cdot (q_0 u_1 + q_1 u_0) d\sigma_y + \int_{\Gamma} \nu \cdot (q_0 u_1 + q_1 u_0) d\sigma_y = 0,$$

and
$$\int_{\mathcal{P}} \nabla_y \cdot \{D(\nabla_x u_1 + \nabla_y u_2)\} dy$$

$$= \int_{\partial Y} D\nu \cdot (\nabla_x u_1 + \nabla_y u_2) d\sigma_y + \int_{\Gamma} D\nu \cdot (\nabla_x u_1 + \nabla_y u_2) d\sigma_y$$

$$= 0 - \int_{\Gamma} \partial_t v_0 d\sigma_y = - \int_{\Gamma} Da(r_S - r_I) d\sigma_y = -|\Gamma| Da(r_S - r_I).$$

Rem: For simplicity, here we assumed that the initial conditions for u and v are ε -independent; this implies the same for u_0 and v_0 .

The Darcy-scale model, non-equilibrium kinetics

Using the cell problems, $U := u_0$, $V := v_0$, q and p solve the effective model

$$\left\{ \begin{array}{ll} q = -\frac{1}{\mu} K \nabla p & \text{in } \Omega, \\ \nabla \cdot q = 0 & \text{in } \Omega, \\ \partial_t U + \frac{|\Gamma|}{|\mathcal{P}|} Da(r_S(U, V) - r_I(U, V)) \\ = \nabla \cdot (A^* \nabla U) - \nabla \cdot (qU), & \text{in } (0, T] \times \Omega, \\ \partial_t V = Da(r_S(U, V) - r_I(U, V)), & \text{in } (0, T] \times \Omega, \end{array} \right.$$

with $A^* \in \mathbb{R}^{d \times d}$, $a_{ij}^* = \int_Y D(\delta_{ij} + \sum_{j=1}^d \partial_{y_j} w^j(y)) dy$

Note: Assuming that the initial data for v is ε dependent, or the grain boundary is heterogeneous ($r = r_\alpha(\frac{x}{\varepsilon}, u, v)$ for $x \in \Gamma^\varepsilon$), then the last equation is y -dependent,

$$\partial_t V = Da(r_S(U, V) - r_I(U, V)), \text{ on } (0, T] \times \Omega \times \Gamma,$$

and

$$\frac{|\Gamma|}{|\mathcal{P}|} Da(r_S(U, V) - r_I(U, V)) \text{ is replaced by } \frac{Da}{|\mathcal{P}|} \int_\Gamma (r_S(U, V) - r_I(U, V)) d\sigma_y.$$

Last remarks

- Homogenization method for deriving upscaled equations approximating problems involving rapid oscillations (parameters, or complex domains).
- Lecture notes (will be updated), slides
- Questions?

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- Next meeting: applications, including scaling and adimensionalisation.
- Many thanks, hope to see you next week, have a haelthy week and a nice weekend!